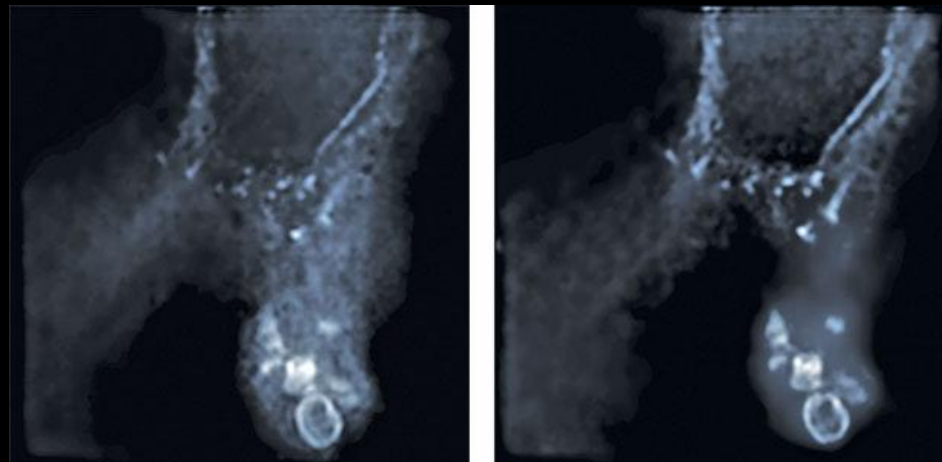
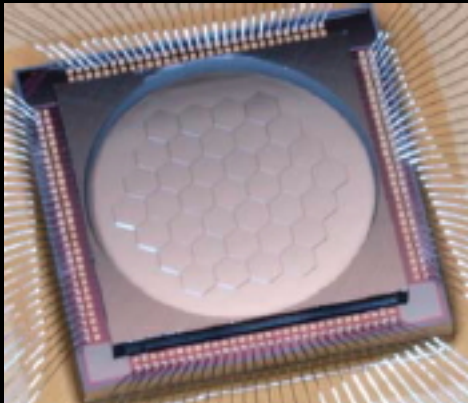


Adaptive Optics

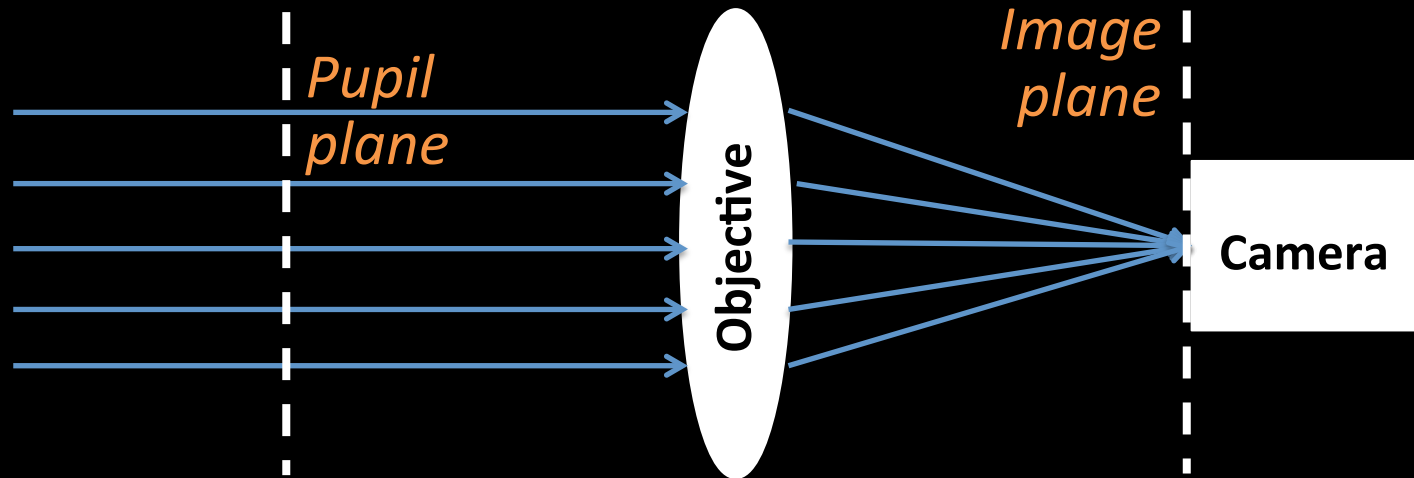
Special Topic in Astrophysics

ASTRON 250 - Fall 2013



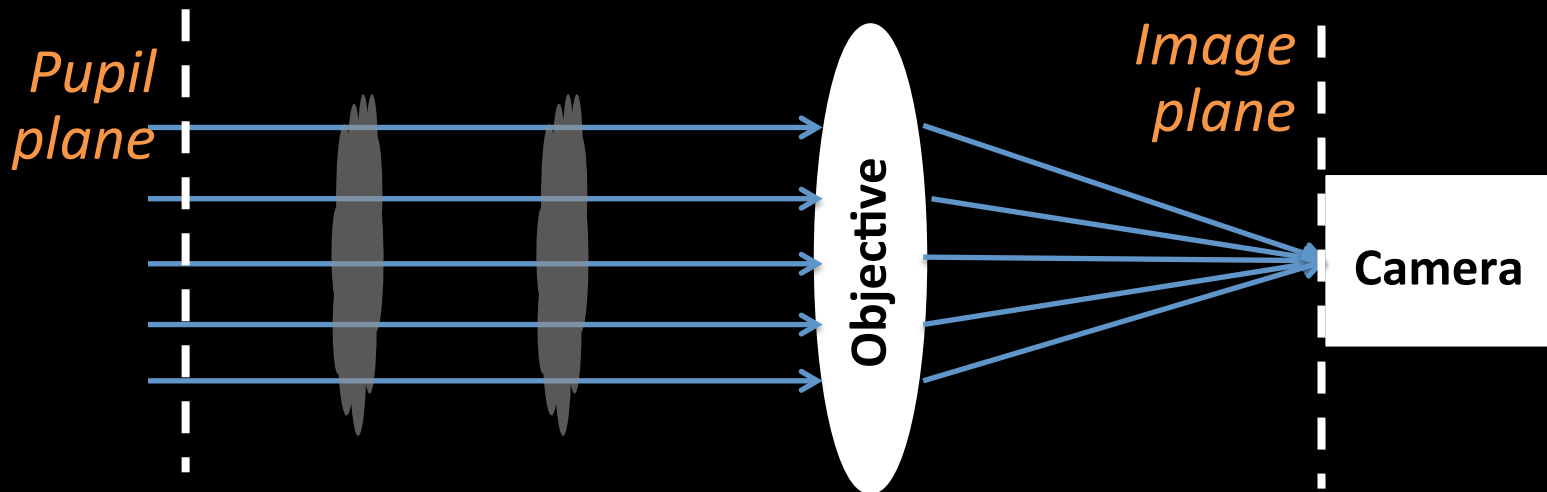
Optical conjugation

- The relationship between the E/B fields in the pupil and image planes is a **Fourier transform**
- Pupil and image planes are **conjugated**



Decomposing the OTF

- $OTF = \prod OTF$ only true if all turbulence phenomena are located in the pupil plane and in the case of small deviations
 - Referred to as the “near field” approximation
- Otherwise, take into account wave propagation



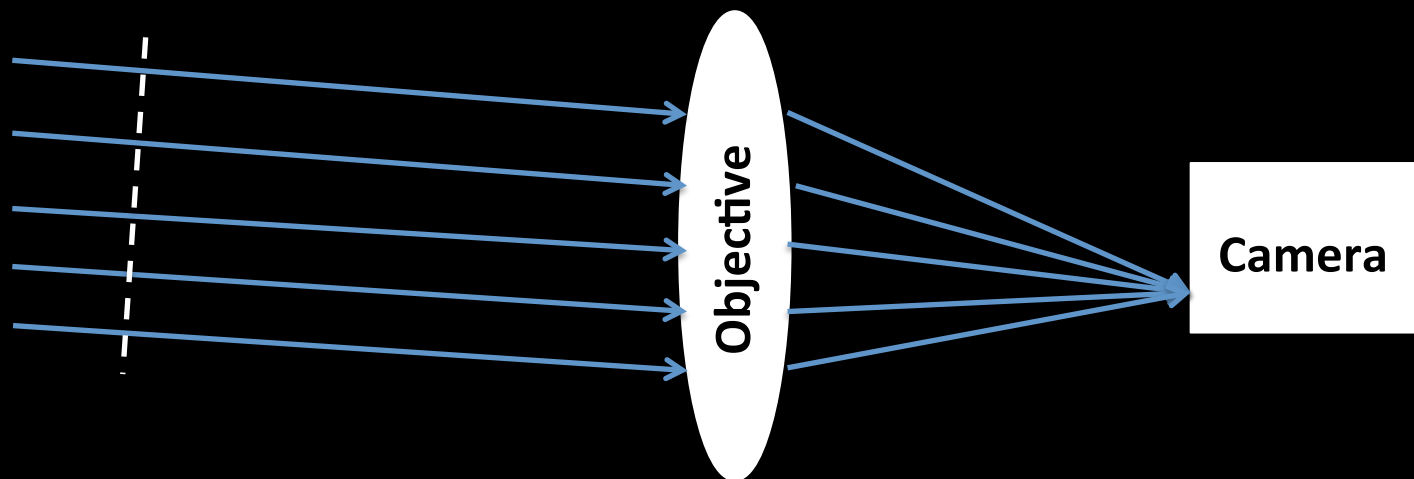
What type of aberrations?

Some simple aberrations

- Defocus
- Chromatic aberration
- Spherical aberration
- Astigmatism, coma
- ...
- Static aberrations can be predicted and corrected for with corrective optics

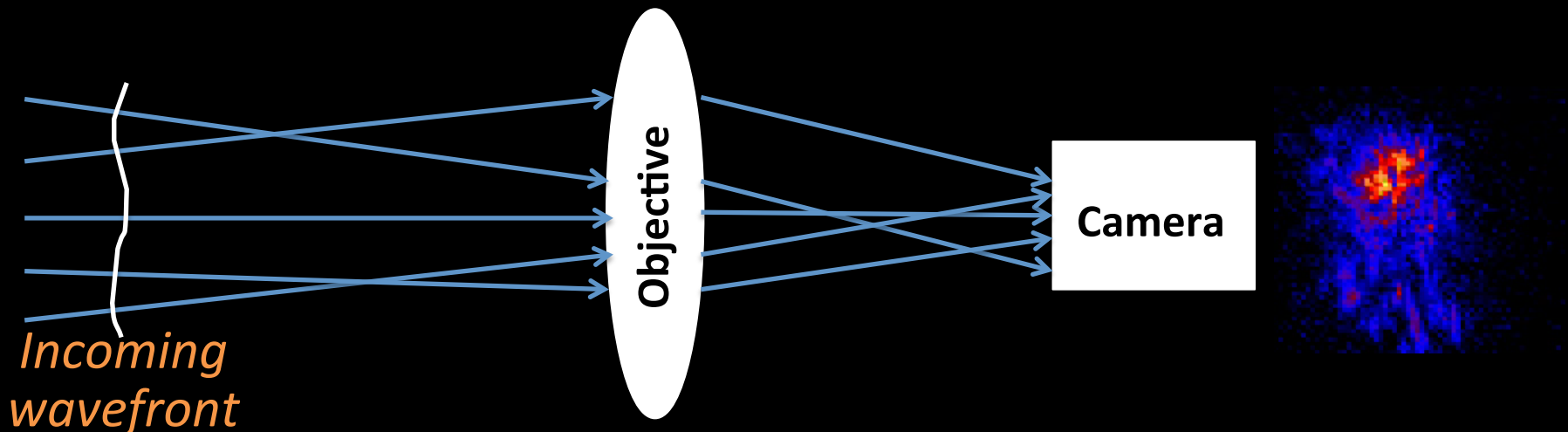
A particularly simple aberration

- Slanted incoming wavefront (**tilt**)
 - Phase in the pupil plane has a linear gradient
 - Results in a **linear displacement** in image plane
 - Image is still **diffraction limited**



More complicated aberrations

- In the general case, wavefronts are corrugated, and each little patch has its own local slope
- Can lead to image blurring or **many displaced diffraction-limited spots** (“speckles”), depending on amplitude of perturbations



Are there “tolerable” perturbations?

- Small enough phase differences do not preclude constructive interference, although it may reduce the resulting intensity
- Larger phase differences are destructive to the signal at the spatial frequency considered
- Phases should remain within $\pi/2$ of each other to ensure good coherence
 - This is Maréchal’s criterion

Describing a corrugated wavefront

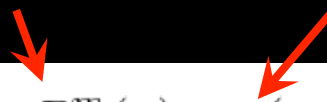
- Any correction relies on the ability to quantify wavefront shape
- Two main options:
 - A **discrete map** of phase delays (and/or slopes) in small patches
 - A **decomposition in a series of analytical functions** that should be selected to be
 - orthogonal to one another (to ensure unicity)
 - “smart” (to minimize the number of terms)

General treatment of aberrations

- A standard method: decomposition in **Zernike polynomials**:

Radial term

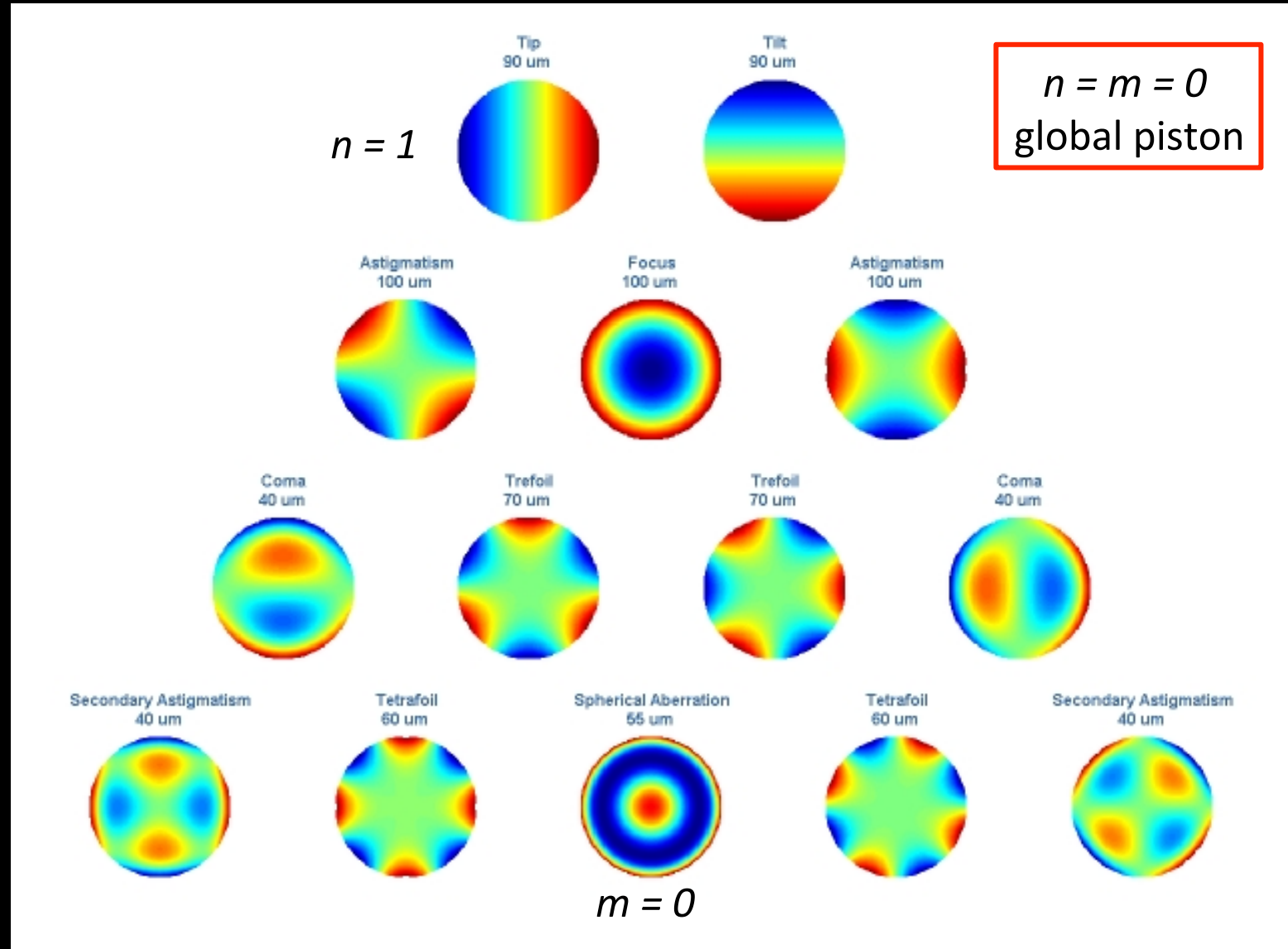
Azimuthal term


$$\begin{aligned} Z_n^m(\rho, \varphi) &= R_n^m(\rho) \cos(m \varphi) && \text{"even" terms} \\ Z_n^{-m}(\rho, \varphi) &= R_n^m(\rho) \sin(m \varphi), && \text{"odd" terms} \\ &&& (n \geq m) \end{aligned}$$

$$R_n^m(\rho) = \sum_{k=0}^{(n-m)/2} \frac{(-1)^k (n-k)!}{k! ((n+m)/2 - k)! ((n-m)/2 - k)!} \rho^{n-2k}$$

if $n-m$ is even, otherwise $R_n^m = 0$

A natural basis to describe aberrations



Know your enemy!

- One needs to understand (and measure) the properties of aberrations to correct for them, or anticipate their effects on the image
 - Ideally precisely (for static aberrations)
 - If they are randomly variable, key statistical properties and power spectrum can be enough

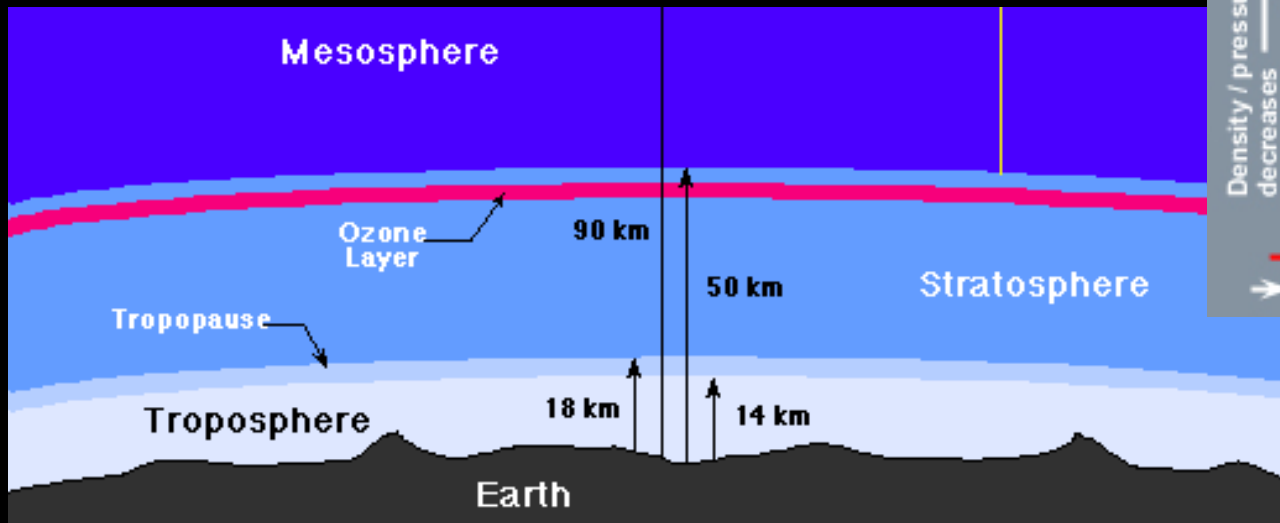
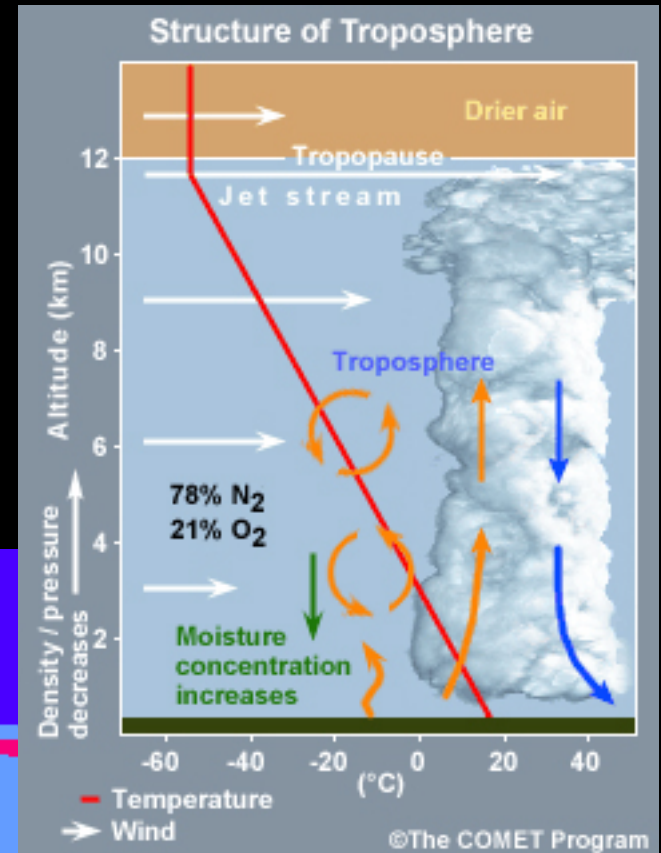
Observing through atmospheric turbulence

- Roddier readings

Observing through atmospheric turbulence

The Earth's atmosphere is characterized by **fluctuations of its index of refraction** due to variations in water vapor content.

Complex motions are also induced by winds (wind shears) and temperature gradients.



Problem set-up

- Atmosphere is constantly changing, so only a **statistical picture** can be drawn
- Atmosphere is described as a number of **independent horizontal layers** whose effects can be added
- In the **frozen flow approximation**, each layer is assumed to be a “phase screen” moving at a constant (wind) velocity

Kolmogorov turbulence

- Assumes an homogeneous atmosphere and isotropic, multi-scale turbulence
- Energy is injected at large scales and transferred to increasingly small scales
 - Inner scale: viscous dissipation (\sim mm-cm)
 - Outer scale: the largest eddies (\sim 10 m)
- Most of the energy is carried by the largest scales: $E(k) \propto k^{-5/3}$

Effect of a single turbulence layer

- Within this model, the structure function (related to the covariance) of the refraction index can be predicted: $D_n(\delta x) = C_n^2 \delta x^{2/3}$
- Refraction index fluctuations result in phase fluctuations, since $d\varphi(x) = 2\pi/\lambda n(x) dz$
- Near-field: phase fluctuations have the same properties at height h and on the ground
 - 2nd order effect: **scintillation** due to diffraction *if the layer is high enough* (and small telescope)

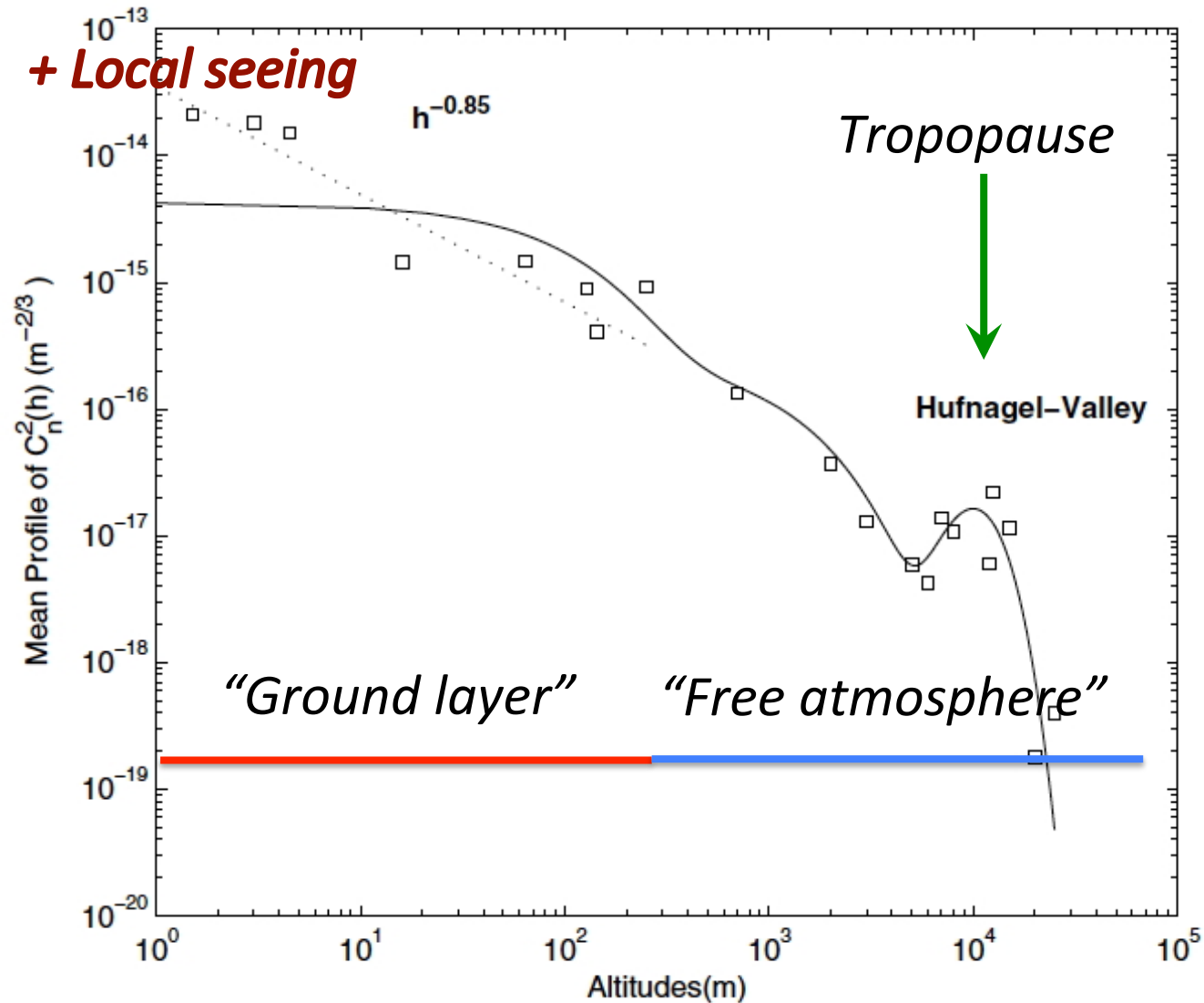
Effect of a single turbulence layer

- For large aperture radii, central limit theorem leads to Gaussian statistics for the phase coherence function (measures the ability of two sub-beams to interfere constructively)
 - $B_{\varphi} = \exp(-1/2 D_{\varphi})$
- After some algebraic manipulations
 - $D_{\varphi}(\delta x) = 2.91 k^2 C_n^2 \Delta h \delta x^{5/3}$
 - $k = 2\pi / \lambda$
 - $\Delta h = \text{layer thickness}$

Adding all layers

- Simply assume **independent horizontal layers**, each one adding a separate phase shift at each position of the wavefront
- Previous equation is modified to becomes a **continuous integral over z**
- All necessary knowledge is encapsulated in the vertical profile of turbulence
 - so-called **C_n^2 profile**

Breaking down the layers

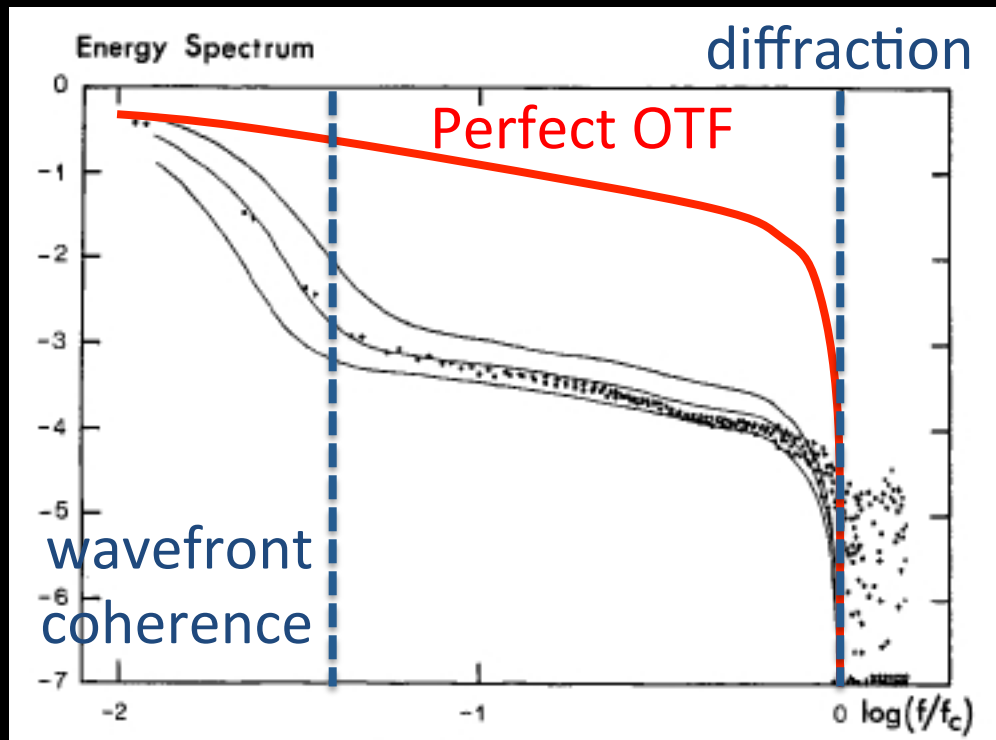


Dealing with multiple layers

- Default is to assume that phase screen are conjugated to primary mirror (near-field)
- However, this is not perfect since the atmosphere is not at infinity
- In case a of a dominant layer, better to work in its conjugated plane rather than in the pupil plane

Effect on OTF

- A full spectrum, attenuating the OTF at all frequencies except the shortest ones (where the field remains correlated)



Fried's parameter

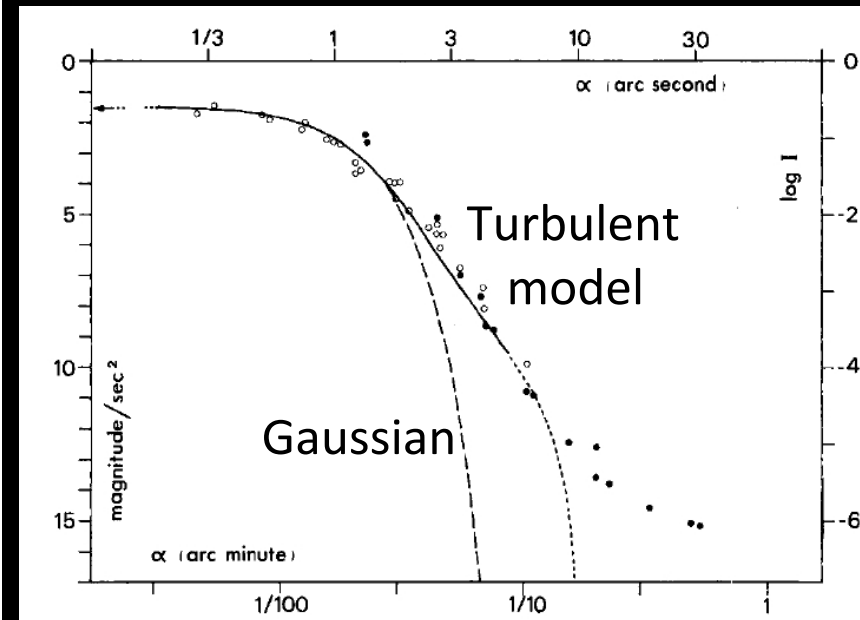
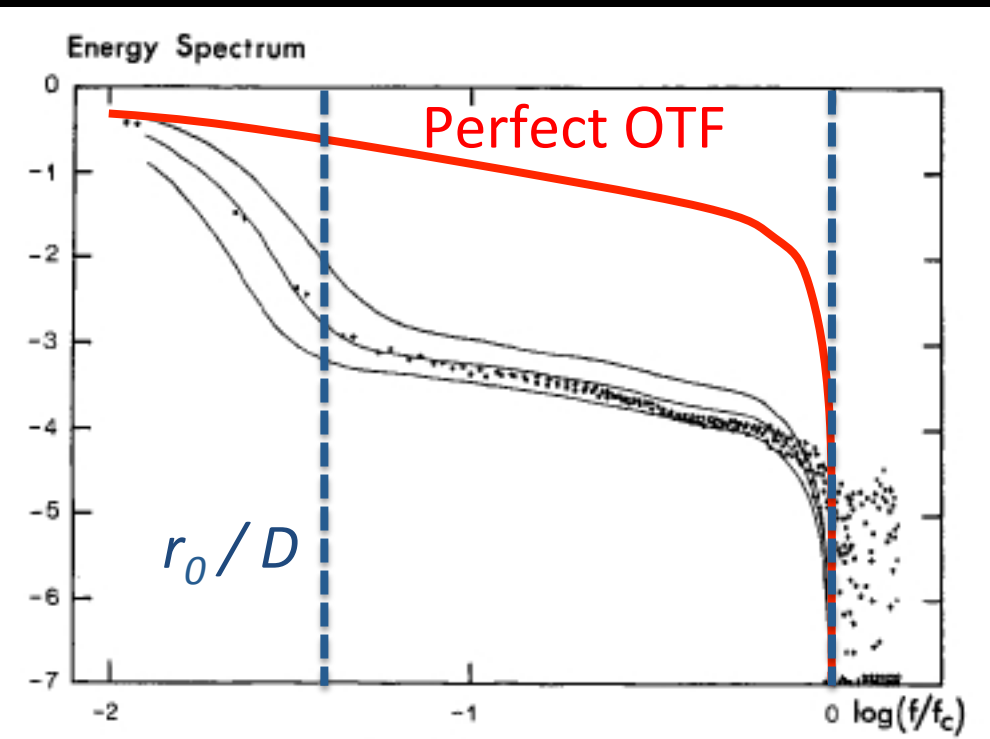
- A natural length scale: $D\varphi(\delta x) = 6.88 (\delta x/r_0)^{5/3}$

$$r_0 = \left[0.423 k^2 (\cos \gamma)^{-1} \int C_N^2(h) dh \right]^{-3/5}$$

- Length scale over which phase correlation is $rms(\varphi) \leq 1\text{rad}$ (remember Maréchal's criterion!)
- r_0 scales as
 - $\lambda^{6/5}$
 - $(\cos \gamma)^{3/5}$ [γ = zenithal angle (sometimes z)]
 - The integral of C_n^2

Fried's parameter

- r_0 drives the “seeing”
- A point source has a quasi-Gaussian profile



Telescope size and r_0

- **Small telescopes:** $D < r_0$
 - Yield diffraction limited images
 - **Still subject to image jitter** (tip-tilt does not decrease wavefront coherence!)
- **Large telescopes:** $D > r_0$
 - Effect of turbulent decorrelation are important
 - Image shape depends on D/r_0 ratio and on **time variability vs exposure time**

Short and long exposure image

- **Short exposure** (“frozen turbulence”) contains information out to the highest frequency
 - It is “diffraction-limited”
 - Independent parts of wavefront form images in different locations: speckle pattern
- **Long exposures** is sum of random short exposures, with constant phase scrambling
 - All high-frequency info is lost, image is completely “turbulence-limited”

Temporal correlation

- Frozen flow + single “effective” wind speed (Taylor approximation)
- Timescale over which the phase is decorrelated ($rms \varphi = 1$ rad) is the $\tau_0 = 0.314 r_0 / v$
- Associated frequency is Greenwood frequency
- Any correction needs to be apply at a higher frequency to be successful

Characterizing an aberrated image

- In the image plane, one can measure the PSF:
 - FWHM, detailed image profile, encircled energy ...
- A common metric: the **Strehl ratio**
 - In an image: **ratio of peak intensity** to that of diffraction-limited image
 - From the OTF: **ratio of integral of the OTF** with and without aberration
 - Equivalent for diffraction-limited images ($S > 0.1$)

Characterizing an aberrated image

$S \sim 60\%$

a

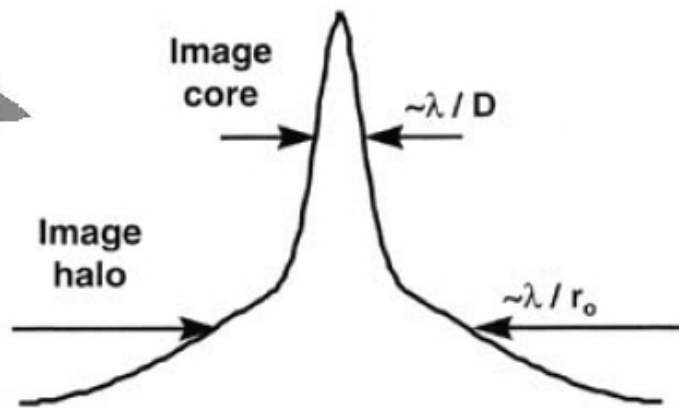
b

$S \sim 20\%$

c

d

$S \sim 5\%$



Quantifying atmospheric turbulence

- Fried & Mevers (1974)
- Racine et al. (1991)
- Tokovinin et al. (2003)

Quantifying atmospheric turbulence

- Estimating r_0
 - Measuring C_n^2
 - Fitting the image of a star
 - Measuring (differential) image motion
 - Measuring (differential) scintillation
- Estimating τ_0
 - Assume some effective wind speed
 - Measure temporal coherence of image motion and/or scintillation

Quantifying atmospheric turbulence

- Typical numbers:
 - $r_0 = 5\text{-}10\text{ cm}$
 - $V = 5\text{-}10\text{ m/s}$
 - $\tau_0 = 1\text{-}5\text{ ms}$
- Variations on all timescales (min to yr)
- No simple connection to weather
- Relative contributions to seeing vary over time

Next week's readings

- Law et al. (2006)
- Labeyrie (1970)
- Babcock (1953)
- Buffington et al. (1977a,b)
- McCall et al. (1977)
- Hinkley et al. (2007)