

SPATIALLY RESOLVED SPECTROSCOPY AND CORONAGRAPHIC IMAGING OF THE TW HYDRAE CIRCUMSTELLAR DISK

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ABSTRACT

We present the first spatially resolved spectrum of scattered light from the TW Hydrae protoplanetary disk. This nearly face-on disk is optically thick, surrounding a classical T Tauri star in the nearby 10 Myr old TW Hya association. The spectrum was taken with the *Hubble Space Telescope* (*HST*) STIS CCD, providing resolution $R \sim 360$ over the wavelength range 5250–10300 Å. Spatially resolved spectroscopy of circumstellar disks is difficult because of the high contrast ratio between the bright star and faint disk. Our novel observations provide optical spectra of scattered light from the disk between 40 and 155 AU from the star. The scattered light has the same color as the star (gray scattering) at all radii except the innermost region. This likely indicates that the scattering dust grains are larger than about 1 μm all the way out to large radii. From the spectroscopic data, we also obtained radial profiles of the integrated disk brightness at two position angles, over almost the same region as previously observed in *HST* WFPC2 and NICMOS coronagraphic images (35 to 173 AU from the star). The profiles have the same shape as the earlier ones, but show a small azimuthal asymmetry in the disk not previously noted. Our STIS broadband coronagraphic images of TW Hya confirm the reality of this asymmetry, and show that the disk surface brightness inside 140 AU has a sinusoidal dependence on azimuthal angle. The maximum brightness occurs at a position angle of $233.6 \pm 5.7^\circ$ east of north. This might be caused by the combination of forward scattering and an increase in inclination in the inner region of the disk, suggesting that the TW Hya disk has a warp like that seen in the β Pictoris debris disk.

Subject headings: planetary systems: protoplanetary disks — stars: individual (TW Hydrae) — stars: pre-main-sequence

1. INTRODUCTION

TW Hydrae is a classical T Tauri star, indicating that gas and dust from a circumstellar (CS) disk are accreting onto the young star (spectral type K7 Ve, distance = $56.4^{+8.1}_{-6.2}$ pc). This disk has now been imaged at visible, near-IR, and millimeter wavelengths (Krist et al. 2000; Weinberger et al. 2002; Wilner et al. 2000). The disk inclination is 7° from face-on (Qi et al. 2004), and no azimuthal asymmetries have been previously reported. A large fraction of the stellar light is reprocessed by the CS dust to far-IR wavelengths ($L_{\text{IR}}/L_* = 0.25$), indicating that the disk is optically thick (Adams et al. 1987). These characteristics suggest that the TW Hya disk is relatively unevolved for its age (~ 10 Myr; Webb et al. 1999).

However, while $\sim 1 \mu\text{m}$ grains in the surface layers of the disk can produce the mid-IR spectrum of the disk, very large (millimeter- to centimeter-sized) grains are needed to explain the millimeter-wavelength spectral energy distribution (Weinberger et al. 2002). The relative lack of near-IR excess flux from the disk indicates that the inner 4 AU have been partially cleared of material, possibly by formation of a giant planet (Calvet et al. 2002). The presence of crystalline silicates in the disk, which are seen in solar system meteorites and comets but not in the interstellar medium, also suggests that the growth of planetary material has begun (Weinberger et al. 2002; Uchida et al. 2004).

Previous broadband photometry indicated that the disk scattering was wavelength independent and that the grain albedo

might be large (Weinberger et al. 2002). This suggests that the disk is composed mostly of icy grains larger than $\sim 1 \mu\text{m}$. However, this conclusion is tentative, since it was based on broadband photometry of the scattered light, rather than actual spectra. This photometry also provided no information about changes in the scattering with location in the disk.

As part of our efforts to understand how the planetary formation process depends on location in the disk, we obtained the first spatially resolved spectrum of scattered light from a protoplanetary disk. Our goals were to study the structure of the TW Hya dust disk, and the size and composition of dust grains as a function of radius from the central star. In addition, we present broadband coronagraphic imaging of TW Hya (see § 7).

2. OBSERVATIONS

Hubble Space Telescope (*HST*) STIS CCD spectra of TW Hya and a point-spread function (PSF) star (CD $-43^\circ 2742 =$ HIP 32939, spectral type M0 V) were obtained on 2002 July 17. All spectra were taken with the G750L grating and GAIN = 1, covering the wavelength range 5250–10,300 Å. The size of the pixels in the spatial direction (y -direction) is $0''.051$.

The sequence of exposures was (1) a point-source target acquisition, (2) a peak-up in a narrow slit ($52'' \times 0''.05$) to accurately determine the position of the star in the dispersion direction (x -direction), (3) a point-source spectrum of TW Hya taken with the $52'' \times 0''.2$ slit, (4) a lamp flat-field image taken with the same slit, and (5) a series of spectra of TW Hya taken

TABLE 1
OBSERVATION LOG

ID	Target	Time (UT)	Exposure Time (s)	Aperture	Comment
Spectroscopy					
2002 July 17:					
o64w57010.....	TW Hya	02:48:46	72	52 × 0.2	Point source
o64w57020.....	Lamp Flat	02:52:08	25	52 × 0.2	
o64w57030.....	TW Hya	04:03:32	198	52 × 0.2F2	Fiducial
o64w57040.....	TW Hya	04:09:02	1380	52 × 0.2F2	Fiducial
o64w57050.....	TW Hya	05:44:00	1380	52 × 0.2F2	Fiducial
o64w57060.....	TW Hya	06:13:18	1221	52 × 0.2F2	Fiducial
o64w57070.....	TW Hya	07:25:00	1380	52 × 0.2F2	Fiducial
o64w58010.....	CD-43 2742	00:38:55	36	52 × 0.2	PSF star, point source
o64w58020.....	Lamp Flat	00:41:41	25	52 × 0.2	
o64w58030.....	CD-43 2742	00:44:41	684	52 × 0.2F2	PSF star, fiducial
o64w58040.....	CD-43 2742	01:00:29	665	52 × 0.2F2	PSF star, fiducial
Imaging					
2000 November 3:					
o64w04010.....	TW Hya	21:01:21	904	WEDGEA1.0	
o64w04020.....	TW Hya	21:19:21	1260.2	WEDGEA1.0	
o64w05010.....	HD 85512	22:29:20	416	WEDGEA1.0	PSF star
o64w05020.....	HD 85512	22:42:08	416	WEDGEA1.0	PSF star
o64w05030.....	HD 85512	22:54:56	416	WEDGEA1.0	PSF star
o64w05040.....	HD 85512	23:07:44	156	WEDGEA1.0	PSF star
2000 December 27:					
o64w06010.....	TW Hya	13:51:48	904	WEDGEA1.0	
o64w06010.....	TW Hya	14:09:48	1260.2	WEDGEA1.0	

NOTE.—All spectroscopic data were taken with the STIS CCD at GAIN = 1, using the G750L grating. All imaging data were taken with the unfiltered (clear) CCD at GAIN = 4.

with the $52'' \times 0''.2$ F2 slit. This slit is the same as the slit used for the point-source spectrum, but with a $0''.86$ wide bar placed over the central star (a fiducial bar). The sequence of exposures was repeated for the PSF star. This observing strategy might be described as “coronagraphic spectroscopy,” which was necessary to reduce the amount of instrumental scattered light and allow detection of the relatively faint disk.

The resolution of the point-source spectra is $R = \lambda/\text{FWHM} \approx 890$ at 7000 \AA (corresponding to $\sim 340 \text{ km s}^{-1}$), while the resolution of the extended-source spectra is $R \sim 360$. The position angle of the slit was 87.4° east of north. A log of observations appears in Table 1.

3. DATA REDUCTION

Since the $52'' \times 0''.2$ F2 slit is an unsupported observing mode, we had to calibrate the data ourselves. We started with the `_crj` exposures produced by the STSDAS IRAF CALSTIS pipeline version 2.14c. Basic data reduction had been performed on these exposures (overscan subtraction, bias subtraction, dark subtraction, simple flat-fielding, and cosmic-ray rejection).

3.1. Alignment of Exposures

Before adding together the `_crj` fiducial exposures, we examined the position of the star in each and found that it did not vary significantly between exposures. Leaving out the first TW Hya fiducial exposure for the moment (o64w57030_crj.fits), the mean x - and y -shifts between exposures were 0.007 and 0.003 pixels. The first TW Hya fiducial exposure was short (only about 14% of the exposure time of the others) and had much lower signal-to-noise ratio, making determination of the star position less precise. Still, the maximum x - and y -shifts of the first exposure

relative to any other were still very small (0.04 pixels in x and 0.09 pixels in y). Therefore, the five TW Hya `_crj` fiducial exposures were added together with no realignment of the images, and the exposure time in the file header changed to the sum of the individual exposure times. The shifts in position of the star between the two PSF fiducial exposures were very small (0.002 pixels in x and 0.005 pixels in y), so they were also combined with no realignment.

3.2. Defringing and Calibration

STIS CCD spectra suffer from fringing at long wavelengths, caused by interference between reflections from the front and back surfaces of the CCD (see Fig. 1). Using our contemporaneous lamp flat images and IRAF tasks provided by STSDAS (`prepspec`, `normspflat`, `mkfringeflat`, and `defringe`), the fringing was removed from the point-source and fiducial data of both the TW Hya and the PSF star. For the fiducial data, the defringing was performed on the combined images, since the higher signal-to-noise ratio allows better defringing than is possible on the separate exposures.

After defringing, we were left with combined two-dimensional spectra in counts per pixel. These were wavelength calibrated using the STSDAS IRAF task `wavecal`. The point-source data were flux calibrated using the STSDAS task `x1d`, which extracts one-dimensional spectra from the images. The fiducial data were flux calibrated with the STSDAS task `x2d`, producing pixels with units of $\text{ergs s}^{-1} \text{ cm}^{-2} \text{ \AA}^{-1} \text{ arcsec}^{-2}$. This task also applies a geometrical distortion correction, producing two-dimensional spectra with wavelength running linearly along the x -axes of the images and distance along the slit running linearly along the y -axes.

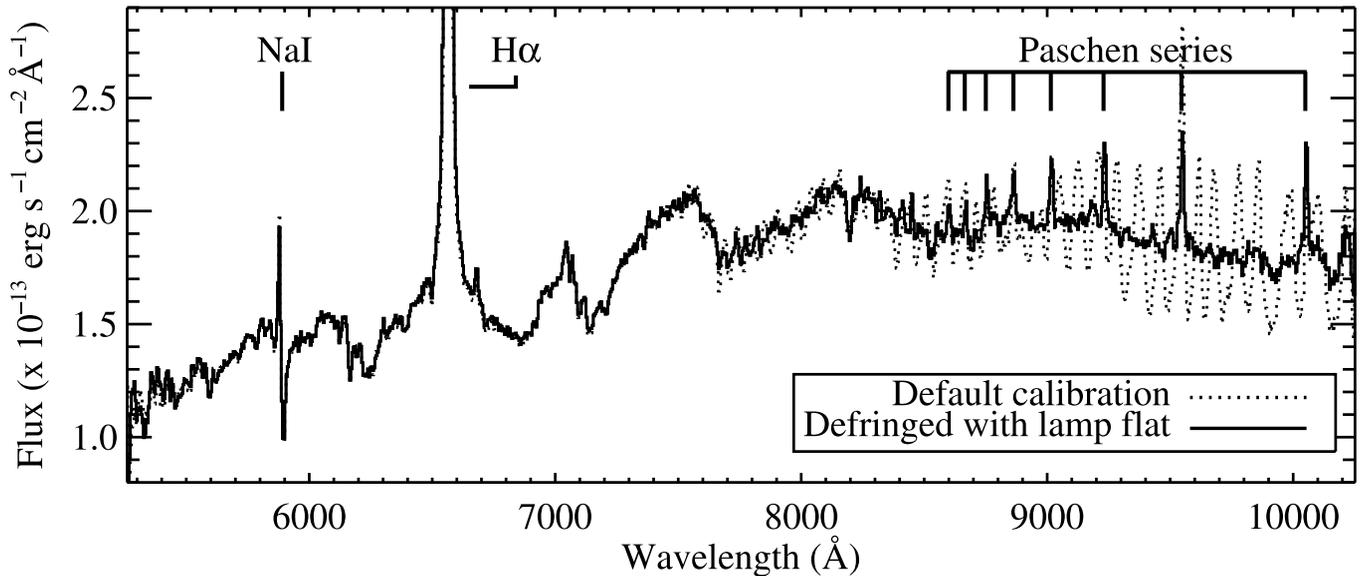


FIG. 1.—Point-source STIS G750L spectrum of TW Hya. The default calibrated data are shown by a dotted line. The defringed data are shown by a solid line. Defringing allows detection of several Paschen series H I emission lines; these lines are caused by accretion of disk gas onto the star, as are the Na I and H α emission lines.

3.3. Correction of Hot and Cold Pixels

We found that the best weekly average dark image used by the CALSTIS pipeline still left a number of hot and cold pixels in the data. We therefore performed a sigma-clipping procedure tailored for two-dimensional spectra on our calibrated fiducial images. The procedure was tailored primarily by using a clip box that mimicked the shape of the data, i.e., 1201 pixels in x (full x -range of the image) by 5 pixels in y . This box worked best since the values along a row of a two-dimensional spectrum do not vary as much as they do in the y -direction; a standard square clip box tended to miss hot and cold pixels in the portion of the image containing the spectrum.

If the brightest (and faintest) pixel in each clip box varied from the median value of its neighboring pixels by more than 3 times the local noise, it was replaced by the median value. The local noise was the standard deviation of the values in the clip box (spectrum noise) added in quadrature to the standard deviation of the column containing the brightest or faintest pixel (background noise). The background noise was determined in this way since it had a clear dependence on x -position (the background is worse on the long-wavelength end of the CCD). This process was applied to the data until the number of deviating pixels approached zero. This clipping procedure did not correct unreasonably large numbers of pixels (only 0.16%), remove background noise, or preferentially treat the portion of the image containing the spectrum.

The final calibrated two-dimensional TW Hya fiducial spectrum appears in the top panel of Figure 2. The primary instrumental features of the two-dimensional spectrum are a smooth halo of scattered light and linear streaks caused by dispersing the Airy rings of the telescope PSF. These streaks angle away from the central star from left to right across the image, since the size of an Airy ring increases with wavelength (see STIS Instrument Handbook, ver. 7.0, § 13.7.3 for an explanation).¹

3.4. PSF Subtraction

The next step in the reduction of the TW Hya data was subtraction of the stellar light from the fiducial image, in order

to isolate light reflected from the CS disk. The PSF star was chosen to match the spectral type of TW Hya as closely as possible, but it is not an exact match. To correct for the difference in overall brightness and spectral mismatch between TW Hya and the PSF star, we made use of the point-source spectra of both stars. Each column of the PSF fiducial image was scaled by the ratio of the TW Hya and PSF star point-source spectra at that wavelength, to produce a synthetic PSF image with the same brightness and color as TW Hya. The statistical errors of the point-source spectra were propagated into the error array of the synthetic PSF image.

The synthetic PSF image was then subtracted from the TW Hya fiducial image; the result is shown in the bottom panel of Figure 2. The remaining continuum flux seen is light reflected from the CS dust disk. At 1'' from the star, the surface brightness of the disk is slightly greater than the surface brightness of the light coming from the star ($S_{\text{disk}} = 1.16 \times S_{\text{star}}$), demonstrating the excellent suppression of scattered stellar light by the fiducial bar.

The disk surface brightness at H α is somewhat greater than it should be, since the emission line in the TW Hya point-source spectrum was saturated, while the line in the fiducial exposures was not. This resulted in underestimated PSF scaling ratios at the saturated wavelengths. We were not able to correct the saturated pixels since the nonlinearity beyond saturation of the STIS CCD at GAIN = 1 has not been characterized (Gilliland et al. 1999). The large-scale factor at H α (caused by the fact that the PSF star does not show the emission line) also produced a stripe of increased noise in the PSF-subtracted image. The H α wavelength region will be ignored in the disk spectra discussed in § 6. A further discussion of this problem appears in § 5.2.

4. SOURCES OF SYSTEMATIC ERROR

The three major factors limiting the quality of a PSF subtraction are (1) a change in the shape of the telescope PSF, (2) an offset in position between the target star and the PSF star, and (3) a color mismatch between the target star and the PSF star. The first factor is caused primarily by thermal flexure of the telescope; there is nothing that can be done postobservation to

¹ See <http://www.stsci.edu/hst/stis/documents/>.

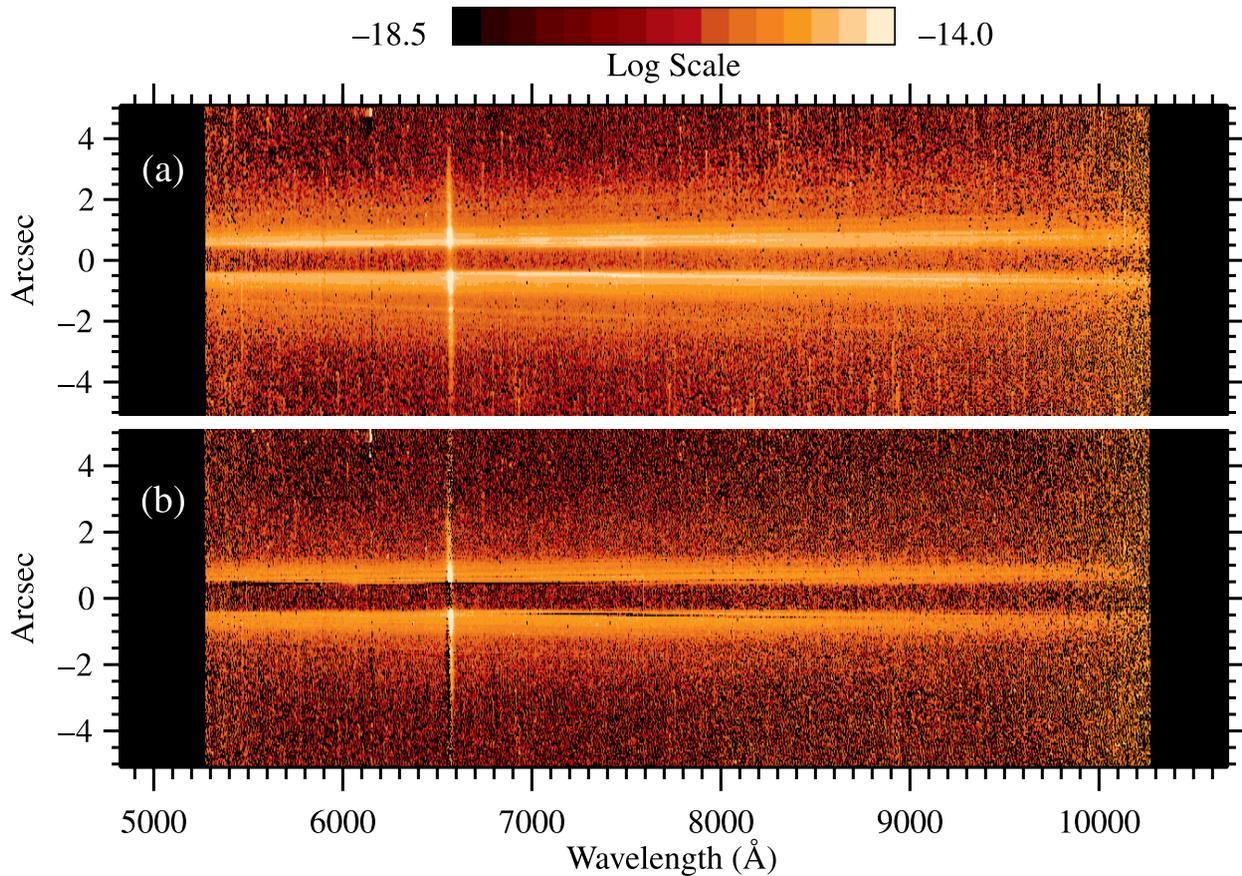


FIG. 2.—STIS G750L fiducial spectra of TW Hya. The dispersion direction is along the x -axis, and the spatial direction is along the y -axis. The figure shows a logarithmic scaling of surface brightness in units of $\text{ergs s}^{-1} \text{cm}^{-2} \text{Å}^{-1} \text{arcsec}^{-2}$. The central star is blocked by the $0''.86$ wide fiducial bar. (a) Total combined TW Hya two-dimensional spectrum. The stellar $\text{H}\alpha$ emission line is visible at 6560 Å . Linear streaks caused by dispersal of the Airy rings of the telescope PSF can be seen angling away from the fiducial bar from left to right across the image. (b) PSF-subtracted two-dimensional spectrum of TW Hya. The light seen is scattered by the CS dust disk. The surface brightness at $\text{H}\alpha$ is somewhat overestimated; see §§ 3.4 and 5.2 for discussion of this issue.

correct this. The only way to minimize the effect is to choose a PSF star that is near the target star on the sky, so that the telescope is oriented in the same way relative to the Sun, and to observe the target and PSF stars close together in time, which we did (observed in consecutive orbits). In §§ 5 and 6, we will discuss the artifacts introduced by this effect. Here we discuss our pre- and postobservation efforts to minimize the other two effects and characterize the systematic errors introduced.

4.1. Misalignment of TW Hya and the PSF

An offset between the positions of TW Hya and the PSF star in the fiducial images will obviously result in a degradation of the quality of the PSF subtraction. To ensure that any offset was as small as possible, a peak-up along the dispersion direction was performed after acquisition of each star. This procedure centers the star in the science slit with an accuracy of 5% of the peak-up slit width, which in our case corresponds to $0.049 x$ -pixels (Brown et al. 2002). Such a small offset does not significantly affect the quality of our PSF subtraction. However, we did not perform a peak-up along the spatial dimension to center the star in the y -direction. A STIS point-source target acquisition typically centers the star to within $0''.01 = 0.2$ pixels (Brown et al. 2002). An offset between TW Hya and the PSF star of a few tenths of a pixel would introduce significant systematic error.

We attempted to measure the offset between the positions of TW Hya and the PSF star in the fiducial spectra several ways. Unfortunately, alignment of the two images appears to be domi-

nated by small-scale structure, and different techniques gave different offsets. As we show in § 5, an offset of $0.0 y$ -pixels produces a PSF-subtracted two-dimensional spectrum with similar total disk fluxes above and below the fiducial. Since we do not expect the disk to be much brighter on one side than on the other, we chose our best y -offset to be zero. In order to investigate the effect of a small misalignment in the y -direction, we produced PSF-subtracted images in which the PSF star image was shifted relative to the TW Hya fiducial image by $\pm 0.25 y$ -pixels. The changes seen in the radial surface brightness profiles and the extracted disk spectra (discussed in §§ 5 and 6) were used to characterize the systematic error introduced by a misalignment of the TW Hya and PSF stars.

4.2. Incorrect Scaling of the PSF Star to TW Hya

Color mismatch between the target star and the PSF star can be a major source of systematic error during PSF subtraction, since the size of the telescope PSF is a function of wavelength. In the case of TW Hya, a mismatch in total brightness may also occur, since this star is variable in brightness and color on various timescales. The Cousins R - and I -band variability is shown below (Mekkadon 1998).
Over 1.13 yr:

$$R_{\text{mean}} \pm 1 \sigma = 10.101 \pm 0.061$$

$$I_{\text{mean}} \pm 1 \sigma = 9.341 \pm 0.031.$$

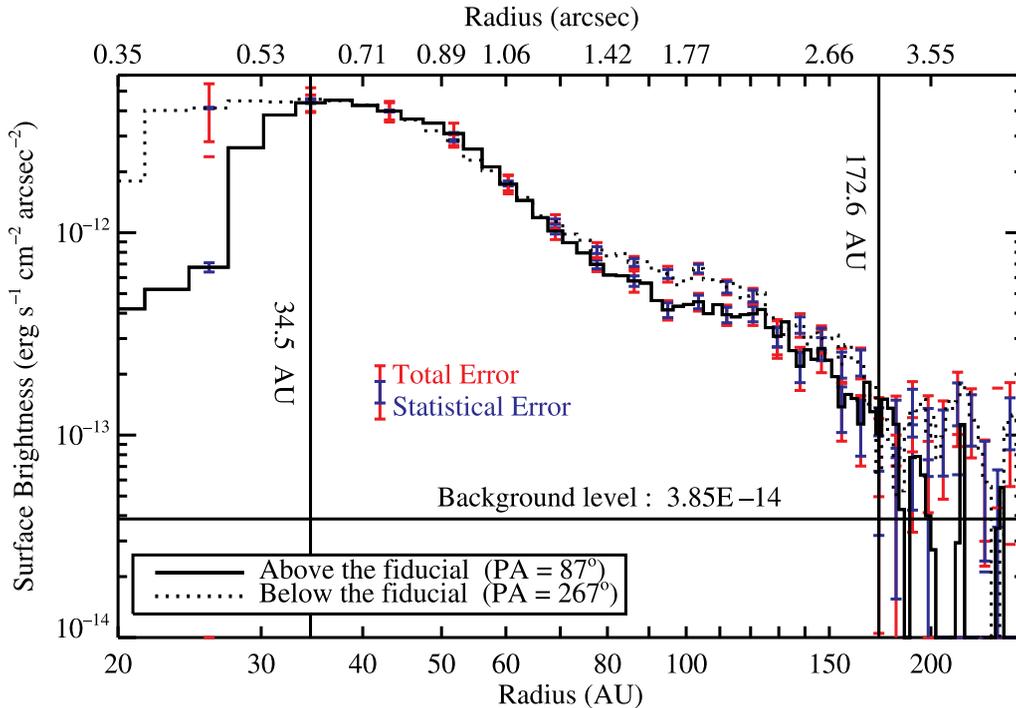


FIG. 3.—Radial surface brightness profiles of the TW Hya disk, integrated over the whole G750L bandpass (5250–10,300 Å). The profile above the fiducial bar (P.A. = 87° east of north) is shown by a solid line; the profile below (P.A. = 267° east of north) is shown by a dotted line. The statistical errors are shown by blue error bars; the total errors, including systematic errors, are shown by red bars. The disk is detected between 34.5 and 172.6 AU. The background level (the median surface brightness outside of 200 AU) is indicated with a horizontal line.

Over 8.03 days:

$$R_{\text{mean}} \pm 1 \sigma = 10.126 \pm 0.048$$

$$I_{\text{mean}} \pm 1 \sigma = 9.351 \pm 0.028.$$

Over 4.12 hr:

$$R_{\text{mean}} \pm 1 \sigma = 10.128 \pm 0.013$$

$$I_{\text{mean}} \pm 1 \sigma = 9.351 \pm 0.009.$$

We addressed the problems of variability and color mismatch by obtaining point-source spectra of both stars before the fiducial observations, so that we would be able to do a near-time, wavelength-dependent scaling of the PSF star to TW Hya. The expected *R*-band variability of TW Hya over the time between the point-source observation and the end of the fiducial observations (approximately two *HST* orbits ≈ 2 hr) is ± 0.0063 , corresponding to a $\pm 0.58\%$ flux variability; the *I*-band variability is smaller. The star should have been stable to within the uncertainty of the STIS flux calibration. We confirmed this by examining the brightness of TW Hya in the five fiducial exposures and found that it was stable at all wavelengths at the 1% level.

Even if the brightness and color of TW Hya were stable over the course of our observations, there still might be an error in the absolute flux calibration of TW Hya or the PSF star. We examined this by adjusting the point-source fluxes by $\pm 1\%$, producing altered ratio arrays and misscaled synthetic images. This changes the overall brightness of the synthetic PSF images and introduces a small color mismatch as well, since the PSF star does not have a flat spectrum. Generally, this introduced a systematic error of $\pm 2\%$ to $\pm 4\%$ in the disk radial profiles and the extracted disk spectra. A few specific pixels in the outer regions of the disk where the signal is low showed larger percent errors.

But as a whole, the error introduced by a reasonable estimate of the possible mismatch between the scaled PSF image and the TW Hya fiducial image is consistently smaller than the statistical uncertainty of the data.

5. DISK RADIAL PROFILES

5.1. Total Profiles

Total disk radial surface brightness profiles were produced by integrating every row of the two-dimensional PSF-subtracted disk spectrum over the whole G750L wavelength band (including the H α emission line). The profiles above and below the fiducial bar are shown in Figure 3. Note that “above the fiducial bar” corresponds to a disk position angle of 87° east of north and “below the fiducial bar” corresponds to P.A. = 267°. The blue error bars are the propagated statistical errors of the data, while the red error bars include the systematic error introduced by misalignment of TW Hya and the PSF image (as discussed in § 4.1). The statistical errors exceed the systematic errors outside of $1''.02 = 57.5$ AU. The inner radius at which we claim detection of the disk (34.5 AU) is the innermost radius at which the profile brightnesses above and below the fiducial are within 1σ of each other. The outermost radius (172.6 AU) is the point at which the profiles begin to be within 1σ of the background level (the median surface brightness outside of 200 AU).

The total disk fluxes are similar above and below the fiducial: $F_{\text{above}} = (6.07 \pm 0.17) \times 10^{-13}$ ergs s $^{-1}$ cm $^{-2}$ and $F_{\text{below}} = (6.44 \pm 0.16) \times 10^{-13}$ ergs s $^{-1}$ cm $^{-2}$. However, the disk is noticeably brighter below the fiducial between $r = 78$ AU and $r = 124$ AU. The difference in surface brightness between the two profiles, integrated over this range of radii, is significant at the 8σ level. The peak difference in surface brightness occurs at 95 AU. This azimuthal asymmetry, the first seen in the TW Hya disk, is $1''.38$ to $2''.19$ from the star, where the statistical uncertainties dominate over the systematic uncertainties and there are no

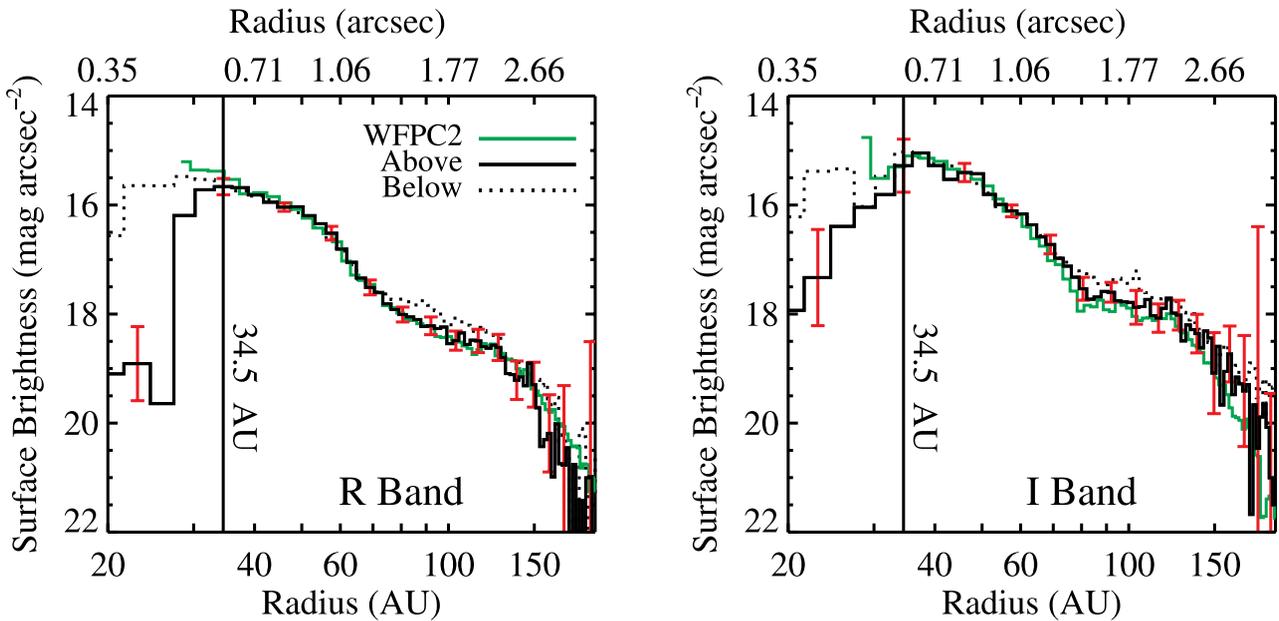


FIG. 4.—*R*- and *I*-band radial surface brightness profiles of the TW Hya disk. The profiles above the fiducial are shown by black solid lines; the profiles below are shown by black dotted lines. The error bars (red) are the total errors, including systematic errors. The *HST* WFPC2 profiles are overplotted with green solid lines. These are the profiles shown in Figure 5 of Krist et al. (2000), converted from WFPC2 F606W and F814W band fluxes per arcsec² to *R*- and *I*-band magnitudes per arcsec². No arbitrary shifting or scaling was applied to the WFPC2 profiles.

visible PSF-subtraction artifacts. In § 7, we discuss our STIS broadband coronagraphic imaging of the disk, which also shows azimuthal asymmetry in the disk brightness.

5.2. *R*- and *I*-Band Profiles

We also produced Cousins *R*- and *I*-band radial surface brightness profiles for comparison to previously published *HST* WFPC2 profiles. The conversion to surface brightness in magnitudes arcsec⁻² was done by applying the following equation to every row of our two-dimensional disk spectrum:

$$S'_{\text{disk}} = m_{\text{Vega}} - 2.5 \log \frac{\int_0^{\infty} q(\lambda) S_{\text{disk}}(\lambda) d\lambda}{\int_0^{\infty} q(\lambda) F_{\text{Vega}}(\lambda) d\lambda}, \quad (1)$$

where S'_{disk} is the disk surface brightness in magnitudes arcsec⁻², m_{Vega} is the magnitude of Vega (0.0), $q(\lambda)$ is the Cousins filter throughput curve from the IRAF SYNPHOT package input tables, $S_{\text{disk}}(\lambda)$ is the disk surface brightness in units of ergs s⁻¹ cm⁻² Å⁻¹ arcsec⁻² (the pixel values in our two-dimensional disk spectrum), and F_{Vega} is the spectrum of Vega from SYNPHOT in units of ergs s⁻¹ cm⁻² Å⁻¹. The resulting *R* and *I*-band radial profiles are shown in Figure 4. The asymmetry seen in the total profile in Figure 3 is seen in both the *R* and *I*-band profiles.

The WFPC2 profiles overplotted in Figure 4 are the profiles shown in Figure 5 of Krist et al. (2000), converted from WFPC2 F606W and F814W band fluxes per arcsec² to *R*- and *I*-band magnitudes per arcsec². Our profiles show the same decrease in slope between 79 and 134 AU, labeled Zone 3 by Krist et al. (2000), and also seen in *HST* NICMOS profiles of the disk (Weinberger et al. 2002). We note that the asymmetry between the profiles above and below the fiducial occurs over almost exactly the same region.

There is extraordinarily good agreement between the absolute brightnesses of our profiles and WFPC2 profiles, which may be somewhat surprising, since the star is variable. We calculated the *R*- and *I*-band magnitudes of the star at the time of

our observations by replacing $S_{\text{disk}}(\lambda)$ in equation (1) with the point-source TW Hya spectrum. The magnitudes of TW Hya were $R = 10.1875 \pm 0.0083$ (including the H α emission line) and $I = 9.4611 \pm 0.0064$, slightly fainter than the mean long-term magnitudes shown in § 4.2. The $\pm 1 \sigma$ errors take into account the 5% absolute photometric accuracy of the STIS flux calibration (Brown et al. 2002).

Saturation of the H α emission line in the TW Hya point-source spectrum will cause our stellar *R*-band magnitude to be somewhat underestimated. We examined this problem by comparing the observed equivalent width of the line to the H α widths seen during a two-year-long TW Hya spectroscopic monitoring project (Alencar & Batalha 2002). We measured a H α equivalent width of $280.5 \pm 6.5 \text{ \AA}$, which is a lower limit to the true equivalent width. This value is very close to the largest equivalent width seen (about $274 \pm 19 \text{ \AA}$; Alencar & Batalha 2002). This shows that although the photosphere of TW Hya was fainter than usual during our observations, the H α line was quite bright; this is not strange, since there is no correlation between the H α line strength and the stellar photospheric flux (Mekkadon 1998; Alencar & Batalha 2002). The large equivalent width of our line also indicates that the H α point-source flux and stellar *R*-band magnitude are probably not greatly underestimated.

The *R*- and *I*-band magnitudes of TW Hya during the WFPC2 observations in Krist et al. (2000) were $R = 9.84 \pm 0.06$ and $I = 9.16 \pm 0.06$, which are 0.35 and 0.30 mag brighter than the star at the time of our observations. The expected 0.35 and 0.30 mag arcsec⁻² differences between our *R*- and *I*-band profiles and the WFPC2 profiles are not seen. We do not know the reason for the close agreement between the profiles.

6. DISK SPECTRA

One-dimensional spectra at increasing radii from the star were extracted from the PSF-subtracted two-dimensional TW Hya disk spectrum. The surface brightness values were integrated over the spatial dimension of the extraction box, then converted to flux units using the *diff2pt* header keyword produced by the

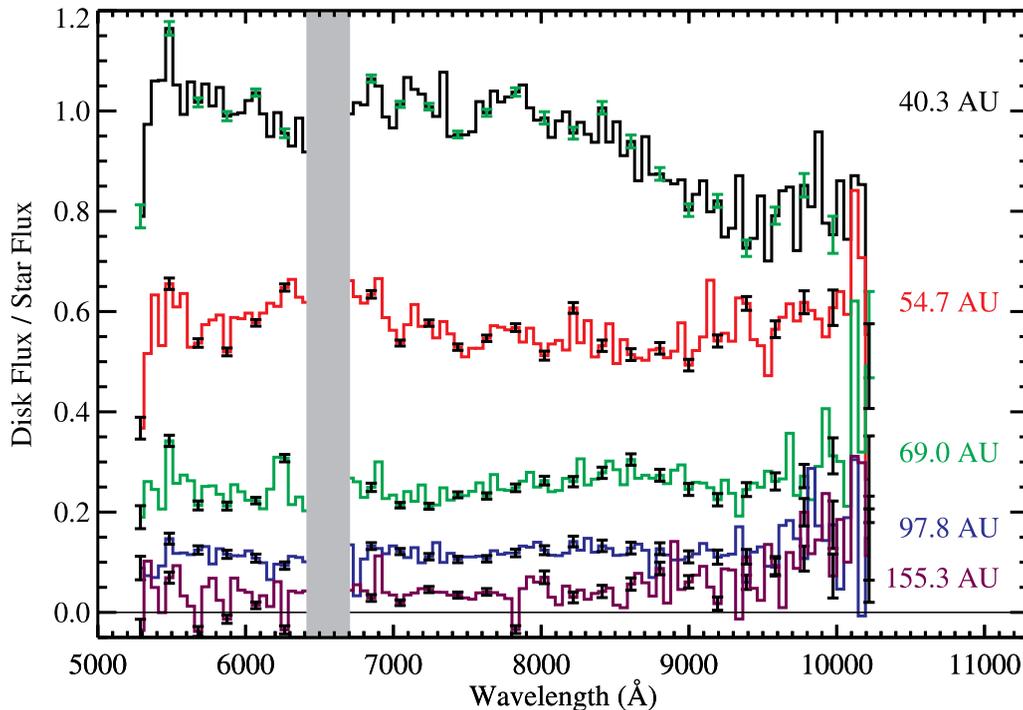


FIG. 5.—Ratio of TW Hya disk spectra to stellar spectrum at increasing radii (referred to in the text as the disk color spectra). The y-axis is the extracted disk flux divided by 2.95×10^{-4} times the TW Hya point-source spectrum. The data have been rebinned by a factor of 10 to increase the signal-to-noise ratio. The region around H α has been blocked out with a gray bar (see § 3.4). The error bars are the $\pm 1 \sigma$ propagated statistical errors.

STSDAS *x2d* tool. This keyword uses a wavelength-averaged aperture throughput value; the extracted spectra have the flux values that are correct in the center of the bandpass but are too high and too low at the ends. This effect is removed when we divide our extracted disk spectra by the point-source stellar spectrum extracted in the same way.

The height of the extraction box was chosen to be 5 pixels ($0''.255$). Spectra extracted using shorter boxes showed peaks and dips due to PSF structure not completely removed by the PSF-subtraction technique. More severe PSF-subtraction residuals near the fiducial bar are the limiting factor on the radius of the innermost extraction box; we avoided the 4 pixels just above and below the fiducial. This centers the innermost spectrum at 40.3 AU, including radii from 34.5 AU ($0''.612$) to 46.0 AU ($0''.816$). Avoiding the 4 pixels nearest the fiducial bar places the innermost pixel of the extraction box at the innermost radius of the disk profiles shown in Figure 3. The PSF-subtraction residuals and their effect on the extracted spectra will be discussed in more detail later in this section.

The disk spectra have the same shape at all radii above and below the fiducial bar, so they were averaged together to increase the signal-to-noise ratio. The extracted disk spectra divided by the arbitrarily scaled TW Hya point-source spectrum (the disk color spectra) are shown in Figure 5, rebinned by a factor of 10 to increase the S/N. Misalignment of TW Hya and the PSF star (discussed in § 4.1) does not significantly affect the disk color spectra. The error bars plotted in Figure 5 are the propagated statistical errors. The S/N per resolution element (two *x*-pixels) of the spectra are 15 at 40.3 AU, 13 at 54.7 AU, 7 at 69.0 AU, 4 at 97.8 AU, and 1 at 155.3 AU. The disk color spectra are flat between 55 and 155 AU, indicating that the scattering function is wavelength independent over the bulk of the disk (neutral or gray scattering).

In the innermost spectrum, the disk-to-star ratio increases at shorter wavelengths; the disk scattering appears significantly

bluer than it does at larger radii. Although we have avoided the worst residuals near the fiducial, we must carefully consider whether this feature is real or due to imperfect PSF subtraction. Unfortunately, the data to accurately quantify the effect of PSF subtraction residuals in STIS fiducial spectra do not exist, since very few observations were taken in this mode. We tried to quantify the effect of residuals in our data in several ways, including analysis of a set of G750L spectra taken with $52'' \times 0''.2$ F2 under calibration program 8844. But these observations were not planned with PSF stability in mind; the best subtraction we could achieve, using the calibration observations closest in time of the two stars most closely matched in spectral type, was obviously much worse than what we achieved for TW Hya.

We are reduced to qualitative arguments about the reality of the blue scattering seen in the spectrum extracted at 40.3 AU. In Figure 6, the innermost region of the PSF-subtracted two-dimensional disk spectrum is shown at the bottom. The worst PSF-subtraction residual is visible as a dark stripe just below the fiducial bar. The spectra extracted at 28.8 and 40.3 AU are shown at the top of the figure; the white bars at the edge of the disk image show the limits of the extraction boxes in the *y*-direction. The spectrum at 28.8 AU is just at the edge of the fiducial bar and includes the worst residual. It can be seen in the extracted spectrum as a roughly 30% dip near 8000 Å.

Since the PSF-subtraction residuals angle away from the fiducial bar from left to right, it might be possible that the worst residual moves into the 40.3 AU extraction box at longer wavelengths. However, the depth of the decrease in the disk color spectrum at long wavelengths is about 30%. A residual that could make such a deep dip in the extracted spectrum should be visible in the two-dimensional spectrum; there is no visible residual in the 40.3 AU extraction box. Also, the disk color spectrum extracted above the fiducial is the same as that extracted below, i.e., bluish. There is no obvious PSF-subtraction residual above the fiducial that could cause this. The blue scattering in the

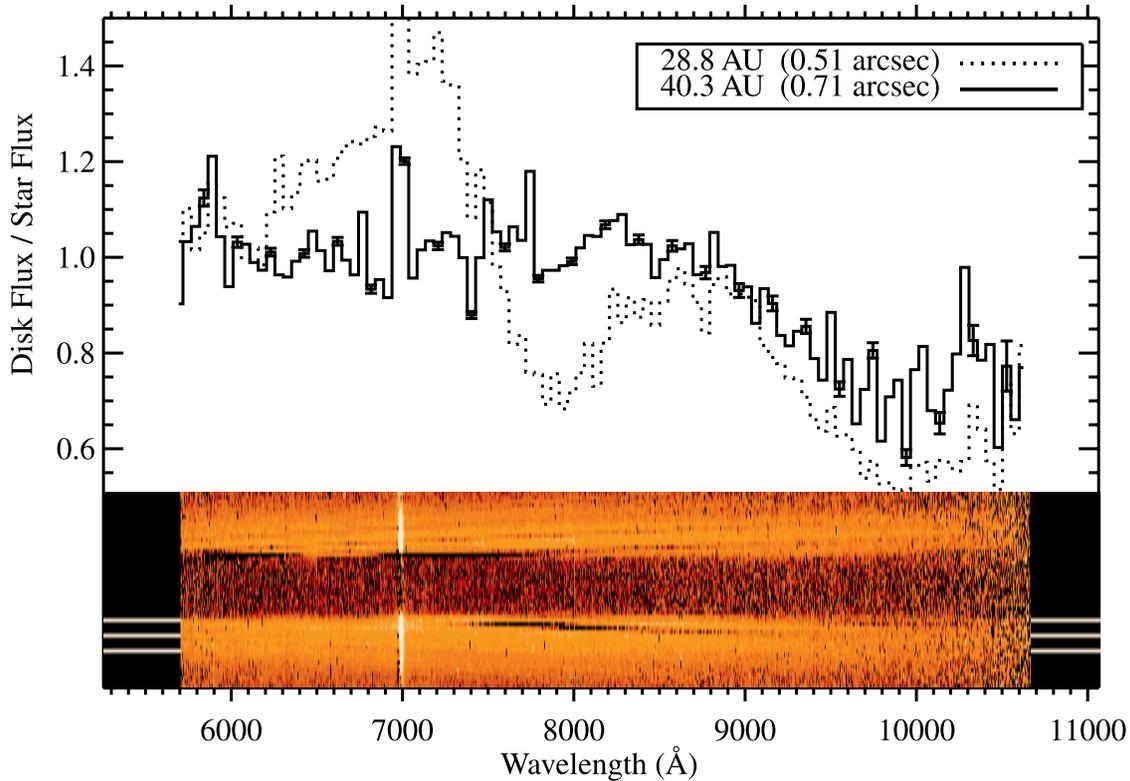


FIG. 6.—Effect of the worst PSF-subtraction residual on the extracted disk color spectra. The PSF-subtracted two-dimensional disk spectrum is shown at the bottom, with the same brightness scaling as in Fig. 2. The spectra extracted at 28.8 AU (*dotted line*) and 40.3 AU (*solid line*) are shown at the top, on the same wavelength scale as the two-dimensional spectrum. The white bars at the edges of the two-dimensional spectrum show the limits of the extraction boxes in the y -direction. The worst PSF-subtraction residual is visible in the two-dimensional spectrum as a dark stripe below the fiducial bar; its effect can be seen in the 28.8 AU spectrum as a dip near 8000 Å.

40.3 AU spectrum is probably real, but we are not able to prove it conclusively.

7. STIS CORONAGRAPHIC IMAGING OF TW HYA

7.1. Observations

STIS CCD images of TW Hya and a different PSF star (HD 85512), were obtained on 2000 November 3 and 2000 December 27. After a standard point-source target acquisition, the stars were observed with the WEDGEA1.0 coronagraphic aperture. This is a $1''$ wide location on the tapered wedge parallel to the CCD charge transfer direction. The STIS CCD is unfiltered, covering essentially the entire visible bandpass. The stars were observed with $GAIN = 4$ to increase the dynamic range of the images. At this gain, the electronics read out the full well of the CCD. Individual frames of 226 s were acquired on TW Hya and 56 s on the PSF, chosen to be longer than the readout time of the CCD for increased efficiency and short enough to avoid saturating the pixels just outside the wedge. The TW Hya field rotated 33° with respect to the STIS CCD between the two visits. The coronagraphic imaging observations are listed in Table 1.

7.2. Data Reduction

CALSTIS performs all of the necessary image reduction, including combination of the cosmic-ray-split exposures. We then combined the repeated exposures, weighting by their integration times, to produce an image in counts s^{-1} . Although bias is supposed to be subtracted in this process, we found that the CCD frames were distinctly nonzero in the portions well away from

the disk. We took the median value of the pixels in the field and subtracted it from the frame to bring the bias to zero.

7.3. PSF Subtraction

As with the spectroscopy, to remove the stellar PSF from the image of the disk, a PSF star must be subtracted. We used the diffraction spikes to align and scale the PSF star separately to the two visits of TW Hya, finding the best subtraction through χ^2 minimization. This process reveals a $(-0.53, 0.00)$ pixel shift between the PSF and TW Hya in the first visit and a $(-0.08, -0.72)$ shift for the second visit. The $\pm 1 \sigma$ error in the alignment is ± 0.027 pixels in x and y . The best-scale factors indicate that the PSF was 3.16 mag brighter than TW Hya in November and 3.08 mag brighter than TW Hya in December. The SIMBAD database gives the V magnitude difference between the two stars as 3.2 mag, so these scales are reasonable. Assuming that the PSF does not vary in brightness, these scalings show that TW Hya was 0.08 mag brighter in November than December. This is within the 2σ variability of TW Hya in the V and R bands on month timescales (Mekkadén 1998).

Next, the December PSF-subtracted image was scaled and registered to the November PSF-subtracted image using the offsets and scale factors measured above. Both images were rotated to put north up and east left. Then, the wedges and 4 pixel wide swaths across the diffraction spikes were masked. The two images were combined by averaging pixels that were unmasked in both images or using pixels that were unmasked in one of the images. The final image is shown in Figure 7. Pixels that were masked in both orientations are blacked out in this image. The slit used in the spectroscopic observations is overlaid in white for illustration.

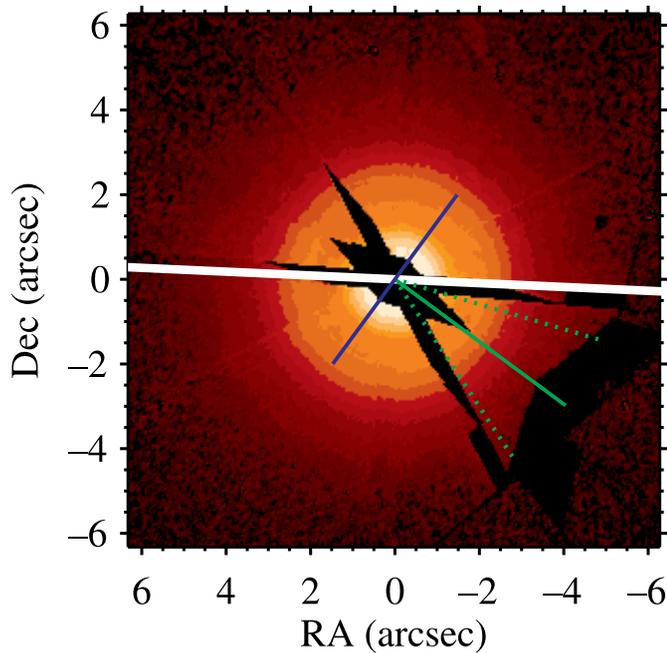


FIG. 7.—STIS CCD coronagraphic image of TW Hya. The image shows a logarithmic scaling of arbitrary surface brightness units. Locations where there are no data because of the coronagraphic wedge or the diffraction spikes are masked out in black. The $0.2''$ wide slit oriented at a P.A. of 87.4° east of north used in the spectroscopic observations is shown in white. The direction of maximum disk brightness (P.A. = 233.6°) found from the sine fitting in Fig. 9 is indicated with a solid green line of length $5.02'' = 283$ AU; this is the outermost radius at which the disk is detected (see Fig. 8). The dotted green lines show P.A. = $233.6^\circ \pm 20^\circ$; pixels between these position angles were included in the maximum brightness radial profile shown in Fig. 8. Our suggested major axis of the inclined inner disk is shown with a solid blue line of length $2.48'' = 140$ AU; this is the outer edge of the inclined inner disk.

7.4. Results

The radial surface brightness profiles from the coronagraphic image are shown in Figure 8. The azimuthally averaged radial profile shows that the disk is detected from 33 AU ($0.59''$) to 283 AU ($5.02''$) in the final coronagraphic image. The inner radius is set by the size of the coronagraphic wedge and the outer radius is the point outside of which the surface brightness is less than 3 times its uncertainty.

Although the disk looks quite circular by visual inspection, a quantitative analysis demonstrates that it is asymmetric. Figure 9 shows the brightness of the disk between 70 and 88 AU as a function of position angle in 20° intervals. The disk brightness peaks at P.A. = $233.6^\circ \pm 4.0^\circ$ and is well fit by a sine function; this error bar is the $\pm 1 \sigma$ statistical error. The direction of maximum brightness is indicated with a green line on the disk image in Figure 7. The ratio of the disk brightnesses in the directions of minimum and maximum brightness is 0.52 ± 0.04 (again, $\pm 1 \sigma$ statistical error). The min/max brightness ratio is relevant to analysis in § 8.1 below.

To estimate the systematic errors introduced by misalignment of the PSF and TW Hya, we offset the PSF image by $\pm 3 \sigma = 0.08$ pixels along the direction of maximum brightness. Offsets in this direction should produce the largest changes to the azimuthal brightness asymmetry. The offset PSF image was subtracted from TW Hya and the sine fit to the brightness versus position angle recalculated. We find that the change to the min/max brightness ratio is only 0.02, which is smaller than the statistical uncertainty (0.04). To estimate the total systematic error in our PSF subtraction, we compared the results from the sine fit to the combined data to the results from fits to the November and December data separately. The difference in the maximum brightness direction between the two visits was 4° , the same as the statistical

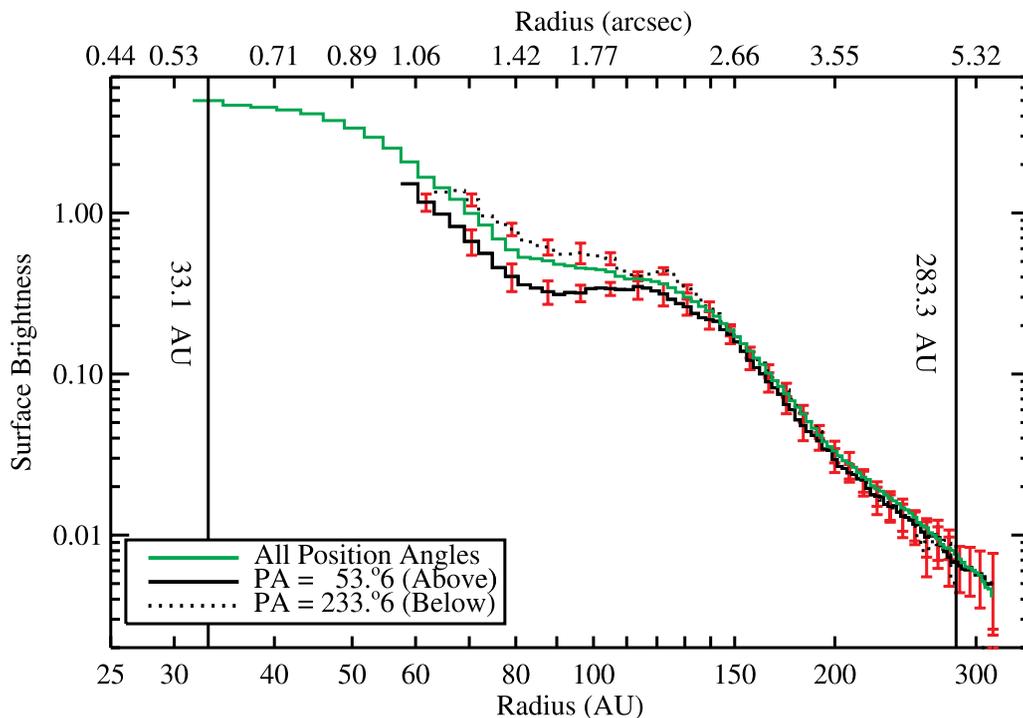


FIG. 8.—Radial surface brightness profiles from the STIS coronagraphic image of TW Hya shown in Fig. 7. Surface brightness in arbitrary units appears on the y-axis, and distance from the star is on the x-axis. The azimuthally averaged profile using all pixels available at each radius is shown with a green line. The disk is detected between 33.1 and 283.3 AU (3σ cutoff). The solid and dotted black lines show the profiles that include only pixels with position angles within $\pm 20^\circ$ of the directions of maximum and minimum brightness found from the sine fitting in Fig. 9. The $\pm 1 \sigma$ statistical error bars on these two curves are shown in red. The position angle of the spectroscopic slit “below the fiducial” (P.A. = 267.4°) is close to the direction of maximum brightness, and the radial profiles from the spectroscopic data shown in Fig. 3 are qualitatively similar to those shown here.

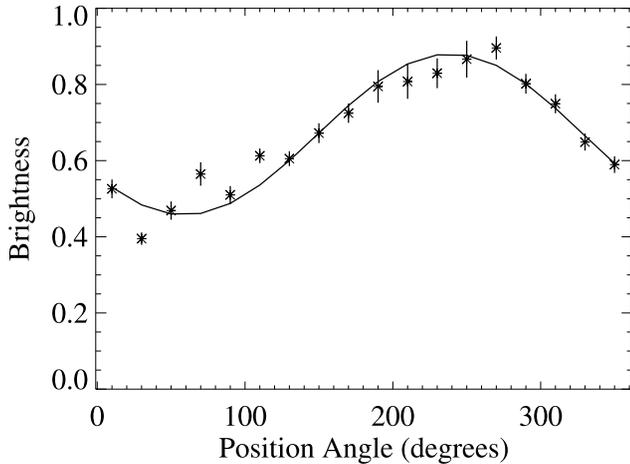


FIG. 9.—Brightness of the TW Hya disk in the STIS coronagraphic image shown in Fig. 7 as a function of position angle. The points show the mean disk brightness between 70 and 88 AU in 20° intervals; the error bars are the $\pm 1 \sigma$ statistical errors. The solid line shows the weighted least-squares sine fit to the data. The directions of maximum and minimum brightness are $233.6 \pm 5.7^\circ$ (including systematic errors) and $53.6 \pm 5.7^\circ$ east of north.

uncertainty of 4° quoted above. Our conservative estimate of the systematic error in the min/max brightness ratio is ± 0.15 . This shows that most of the systematic error is not due to misalignment but is most likely due to changes in the shape of the PSF.

The profiles including only pixels with position angles within $\pm 20^\circ$ of the directions of maximum and minimum brightness are also shown in Figure 8. Unfortunately, the maximum brightness profile (P.A. = 233.6°) only reaches in to 65 AU and the minimum brightness profile (P.A. = 53.6°) only reaches in to 62 AU because of the position of the coronagraphic wedge in both visits. But the profiles clearly show that the disk is significantly brighter on one side between 65 and 140 AU. The difference in surface brightness between the two profiles, summed over this range of radii, is significant at the 10σ level. Outside 140 AU, the disk is symmetric.

The position angle of the slit in the spectroscopic observations “below the fiducial” is only 34° away from the direction of maximum brightness, and comparison of the profiles in Figures 3 and 8 is instructive. The profiles are qualitatively similar; both show that the disk is brighter in the “below the fiducial” direction over some range of radii around $r \sim 100$ AU. However, the coronagraphic profiles show that the disk is significantly asymmetric over a larger range of radii than the spectroscopic profiles do. The reason for this may be the larger error bars of the spectroscopic profiles and/or the arbitrary alignment of the TW Hya and synthetic PSF fiducial images in the y -direction before subtraction, as discussed in § 4.1.

8. DISCUSSION

8.1. Azimuthal Asymmetry in the Disk Radial Profiles

Our spectroscopic and coronagraphic disk radial surface brightness profiles show the first azimuthal asymmetry seen in the TW Hya disk. The coronagraphic profiles show that the disk is brighter on one side than the other between 65 and 140 AU, with the maximum brightness occurring at P.A. = $233.6 \pm 5.7^\circ$ (including the systematic uncertainty). We are not sure exactly how far in toward the star this asymmetry extends, because of the limitation of the wedges in the coronagraphic images and the possible misalignment in the y -direction of the TW Hya and PSF spectroscopic fiducial images.

Forward scattering can cause a sinusoidal brightness asymmetry between the near and far sides of an inclined, optically thick disk (e.g., GG Tau A; McCabe et al. 2002). However, we must explain why the asymmetry is only seen inside 140 AU. The first suggestion is that the grains in the inner disk are more strongly forward scattering. However, the disk color spectra show no change between 55 and 155 AU, above or below the fiducial (see Fig. 5), leading us to believe that the asymmetry is not due to a change in grain composition or size at 140 AU. The second suggestion is that the disk inclination is larger interior to this radius, i.e., that the disk is warped. In this scenario, the direction of maximum brightness is the minor axis of the inclined inner disk, and the P.A. of the major axis is $323.6 \pm 5.7^\circ$. This P.A. roughly agrees with that deduced from submillimeter CO emission observations showing the Keplerian rotation of the disk ($-45^\circ = 315^\circ$; Qi et al. 2004).

If the asymmetry is due to forward scattering, then its magnitude may be used to calculate a lower limit to the inner disk inclination. The minimum to maximum surface brightness ratio from the sine fitting discussed in § 7.4 is $S_{\min}/S_{\max} = 0.52 \pm 0.16$, including the systematic error. Using a Henyey-Greenstein scattering phase function (Henyey & Greenstein 1941), the ratio of the minimum to maximum surface brightness is

$$S_{\min}/S_{\max} = \frac{[1 + g^2 - 2g \cos(90^\circ - i)]^{3/2}}{[1 + g^2 - 2g \cos(90^\circ + i)]^{3/2}}, \quad (2)$$

where g is the dust asymmetry parameter ($g = 1.0$ indicates purely forward scattering, $g = 0$ indicates isotropic scattering) and i is the disk inclination ($i = 0^\circ$ corresponds to face-on). Using the maximum possible value of $g = 1.0$, we find an inner disk inclination of $12.4^{+6.8}_{-5.0}$ deg. This is barely consistent with the inclination of $7^\circ \pm 1^\circ$ found by Qi et al. (2004) from their submillimeter CO data. However, note that the submillimeter data probe a wide range of radii in the disk and give some sort of mean disk inclination.

Using a more plausible value of $g = 0.5$, we find an inclination of $15.6^{+8.7}_{-6.4}$ deg. Our lower limit ($i > 9^\circ$) is significantly greater than the upper limit on the inclination found from isophotal fitting to the TW Hya NICMOS imaging data ($i < 4^\circ$; Weinberger et al. 2002). Our data have 50% better spatial resolution than the NICMOS coronagraphic images. Given that the region over which the disk appears asymmetric is small, i.e., 65–140 AU, or only $1/3$, the STIS images are uniquely suited for measuring the inclination.

A warped inner disk has been observed in the β Pictoris debris disk (see Fig. 8 in Heap et al. 2000). Two possible explanations have been put forward for disk warps (Heap et al. 2000). In the first, warps form through radiation-induced instability in accretion disks surrounding luminous stars (Armitage & Pringle 1997). The problem with this explanation is that the star must have a luminosity greater than about $10 L_\odot$, while the TW Hya luminosity is only about $0.25 L_\odot$ (Webb et al. 1999). In the other explanation, the warp is caused by the gravitational perturbation of a giant planet on an inclined orbit. The presence of a protoplanet in the TW Hya disk has already been suggested to explain the dearth of material within 4 AU of the star (Calvet et al. 2002). However, it is not clear whether a single protoplanet can explain both the central hole and the disk azimuthal brightness asymmetry.

8.2. Gray Scattering over the Bulk of the Disk

The scattering over the bulk of the disk is clearly wavelength independent in the optical, in contrast to the results of Krist et al. (2000), who found the disk to be bluer than the central star.

To determine what this might mean for the properties of the TW Hya dust, we examine the McCabe et al. (2002) Monte Carlo scattering simulation of GG Tau A, a classical T Tauri star binary system surrounded by a massive disk.

The scattered surface brightness at any one angle, S_{scatt} , is proportional to $\omega^n(g_\lambda)$, where ω is the grain albedo and n is the number of scattering events. An increase in n produces a decrease in S_{scatt} . The number of scattering events depends on the dust asymmetry parameter, g_λ , which is a function of wavelength. Note that in the case of TW Hya, the relatively low inclination of the disk indicates that the observed photons have all scattered through almost the same angle ($\approx 90^\circ$).

The wavelength dependence of the scattered light is determined by the combination of the wavelength dependences of ω and g_λ . The wavelength dependence of the albedo tends from blue (Rayleigh scattering; small particles) to neutral (large particles). The dust asymmetry parameter increases with the size parameter $x = 2\pi a/\lambda$, reaching a roughly constant value at $x \gtrsim 4$ (Wolff et al. 1998). Since an increase in g_λ produces an increase in n and a decrease in S_{scatt} , the effect of the dust asymmetry parameter is red (small particles) to neutral (large particles) scattering. Therefore, an optically thick scattering situation may produce any color light. Which parameter dominates depends on the number of scattering events. For example, in a relatively low-density case (low n), the wavelength dependence of S_{scatt} is dominated by ω . The scattered light will be blue to neutral, depending on the particle size. In a high-density case (high n), g_λ will dominate, and the scattered light will be red to neutral.

So there are two ways to get neutral scattering from a disk. In the first, the number of scattering events is such that a blue albedo is exactly canceled by a red dust asymmetry parameter; this implies relatively small scattering particles. In the other, the scattering grains are so large that ω and g_λ are wavelength independent. For the first scenario to produce neutral scattering from the TW Hya disk over such a large range of radii (≈ 100 AU), the number of scattering events would have to be just right and not change much with radius. It seems more likely that the second scenario applies and relatively large grains are present between about 55 and 155 AU. Our disk spectra show that the bulk of the disk neutrally scatters wavelengths between about 5500 and $1 \mu\text{m}$; this indicates grains larger than $x\lambda/2\pi \approx 4(1\mu\text{m})/2\pi \approx 0.6 \mu\text{m}$. Since the disk appears to scatter neutrally all the way into the near-IR ($1.6 \mu\text{m}$; Weinberger et al. 2002), the particle size is likely greater than $1 \mu\text{m}$. These arguments and conclusions will be quantitatively tested in a future paper by application of the McCabe et al. (2002) Monte Carlo scattering code to the TW Hya disk.

8.3. Scattering Change in the Innermost Region of the Disk

We now consider what the apparent change in the wavelength dependence of the scattered light inside 55 AU might mean. From

the above arguments, blue scattering occurs when both the number of scattering events and the grain sizes are relatively small. If we imagine for a moment that the scattering particles are small over the whole disk, the blue scattering at small radii could be caused solely by a decrease in the number of scattering events with decreasing radii. However, since we previously argued that the grains in the outer regions of the disk are large, the change in scattering at small radii must also indicate a decrease in particle size.

The idea that there are more small particles close to the star at first seems counterintuitive, since the coagulation of condensed grains occurs faster at smaller radii. However, the destruction of planetary material also occurs first at smaller radii, as evidenced by central holes in many debris disks. Submicron grains in debris disks, like the 12 Myr old β Pictoris system (Zuckerman et al. 2001), must be produced by collisions among planetesimals (Artymowicz 1988). We therefore suggest that blue scattering at small radii in the TW Hya disk may be evidence that the clearing out of solid planetary material has begun in the inner disk.

9. CONCLUDING REMARKS

These observations demonstrate the usefulness of high-contrast coronagraphic spectroscopy of scattered light from CS disks. We here summarize our principal observational results from this study of the TW Hya protoplanetary disk.

1. Inside 140 AU, the disk surface brightness shows a sinusoidal dependence on azimuthal angle. The maximum brightness is at P.A. = 233.6 ± 5.7 . The origin of this feature requires further study, but we suggest it is caused by an increase in inclination in the inner part of the disk, i.e., a warp.
2. The disk spectra definitively show that the scattering is wavelength independent in the optical between 55 and 155 AU from the star. We argue this indicates that the scattering grains are larger than about $1 \mu\text{m}$ out to large radii.
3. The color of the scattered light from the innermost portion of the disk (34.5–46.0 AU) appears bluer than the outer disk. This result is tentative because of possible systematic errors in PSF subtraction close to the fiducial bar. However, taken at face value, the change in color might indicate a decrease in grain size at small radii.

We thank John Krist for providing us with the WFPC2 radial surface brightness profiles of TW Hya for inclusion in Figure 4. This paper is based on observations made with the NASA/ESA *Hubble Space Telescope*, obtained at the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Inc., under NASA contract NAS 5-26555. These observations are associated with program GO-8624; we thank our co-investigators on this program for useful discussions.

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